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VACUUM PERMEATOR ANALYSIS FOR EXTRACTION OF TRITIUM FROM DCLL BLANKETS

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It is envisioned that tritium will be extracted from DCLL blankets using a vacuum permeator. We derive here an analytical solution for the extraction efficiency of a permeator tube, which is a function of only two dimensionless numbers: one that indicates whether radial transport is limited by the PbLi or by the solid membrane, and another that is the ratio of axial and radial transport times in the PbLi. The permeator efficiency is maximized by decreasing the velocity and tube diameter, and increasing the tube length. This is true regardless of the mass transport correlation used; we review several and find that they differ little, and the choice of correlation is not a source of significant uncertainty here. The PbLi solubility, on the other hand, is a large source of uncertainty, and we identify upper and lower bounds from the literature data. Under the most optimistic assumptions, we find that a ferritic steel permeator operating at 470 °C will need to be about twenty times larger in volume than previous conceptual designs using niobium and operating at higher temperatures.

I. INTRODUCTION

A key safety requirement for any tritium breeding system for a Fusion Nuclear Science Facility (FNSF), demonstration power plant (DEMO), or fusion power plant is minimizing tritium permeation into the reactor building and subsequent release to the environment. The most effective way to meet this requirement is with an efficient tritium extraction system. The Dual-Cooled Lead Lithium (DCLL) breeding blanket,^{1,2} a favored design in US TBM (Ref. 3) and power plant⁴ concepts, is no exception. One significant advantage of the DCLL is that, because of its high temperature and high PbLi flow rate relative to other PbLi blanket designs, it can achieve high thermal efficiencies. However, the higher flow rates also create somewhat different requirements for the tritium extraction system, which must be able to handle the high throughput while maintaining low tritium concentrations in order to reduce overall system permeation losses. For that reason, the compact mass extractor (gas-liquid contactor) envisioned for the Helium-Cooled Lead Lithium (HCLL) design is probably less suitable for the DCLL (Ref. 5).

Instead, it has been envisioned since the US TBM program³ and in subsequent power plant studies⁴ that tritium will be extracted from PbLi using a vacuum permeator. Such a device seeks to exploit the permeation mechanism for extraction by passing the fluid along a high-permeability membrane, with a vacuum maintained on the other side, so as to maximize the concentration gradient across it. The concept has long been in use in gas systems, typically with Pd-Ag membranes.⁶ In PbLi systems, the situation is slightly more complicated. Unlike in gases, resistance to tritium transport radially across the PbLi boundary layer may be significant; this process may be a limiting factor in the efficiency of the device.

As with gas systems, PbLi vacuum permeator models^{4,7} originally sought to use high permeability membrane materials, such as alloys of niobium or tantalum, which could operate at high temperatures (700-730 °C). While excellent for their high permeability, these group 5 metals readily oxidize at high temperatures even at extremely low oxygen partial pressures.⁸ In addition, corrosion will result in some niobium circulating in the PbLi, which may be activated in the blanket and lead to long term waste disposal concerns. Therefore, a permeator employing more conventional materials such as RAFM steel is of considerable interest; it would need to operate at lower temperatures (~470 °C) to prevent excessive corrosion of the RAFM in PbLi (Ref. 9).

In order to better understand the relative importance of different transport phenomena in PbLi permeators, and how their efficiency scales with various design parameters, we derive here a quasi-2D analytical solution for the concentration profile along the axis of a permeator tube and for the overall efficiency of a permeator. In section III, we use this formulation to assess the feasibility of a ferritic steel permeator.

II. ANALYTICAL SOLUTION

II.A. Radial Tritium Transport

As in its original conception,⁷ we envision a vacuum permeator as an array of N cylindrical tubes of inner diameter $d=2r_i$ and outer diameter $d_o=2r_o$ (thickness $r_o -$

r_i) with tritium-laden PbLi flowing in their interiors, and a vacuum maintained on the collective exterior. As in previous analyses,^{3,4,7} we consider three radial transport processes, shown schematically in Figure 1.

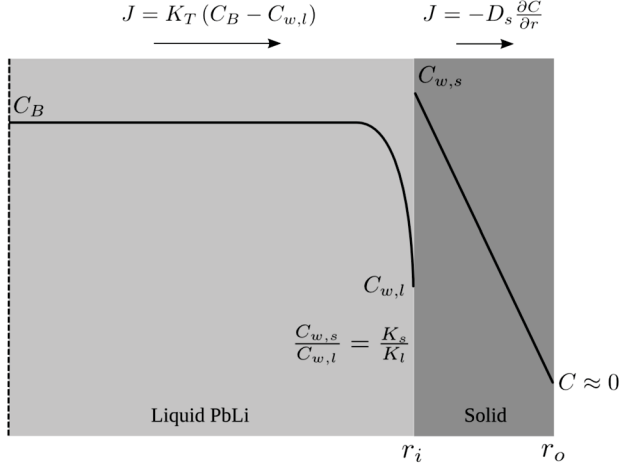


Fig. 1. Schematic of tritium transport processes in the radial direction in a vacuum permeator tube.

The radial flux in the liquid (subscript l) PbLi is determined by the mass transport coefficient, K_T (m/s):

$$J_l = K_T (C_B - C_{w,l}) \quad (1)$$

where C_B and $C_{w,l}$ (mol T/m³) are the tritium concentrations in the bulk and at the wall, respectively. The latter is presumed to be in equilibrium with the solid (subscript s) concentration $C_{w,s}$ at the same interface, in which case they are related by the ratio of their solubilities:

$$\frac{C_{w,s}}{C_{w,l}} = \frac{K_s}{K_l} \quad (2)$$

where K_s and K_l (mol T/(m³·Pa^{1/2})) are the solid and liquid solubilities, respectively. The concentration in the solid is given by the well-known solution to the diffusion equation in cylindrical coordinates (see e.g. Ref 10),

$$C = \frac{C(r_i) \ln(r_o/r) + C(r_o) \ln(r/r_i)}{\ln(r_o/r_i)} \quad (3)$$

Since $C(r_i) = C_{w,s}$ and $C(r_o)$ is assumed to be near zero as a result of the vacuum on the outside of the tubes, Fick's law gives the permeation flux at the interface as

* Note that when employing the resistance analogy, as is done in Refs. 11-12, this discontinuity implies an additional (negative) "contact resistance" of the form $(1/h_c) = ((K_l/K_s)-1) r_i \ln(r_o/r_i)/D_s$.

$$J_s = \frac{D_s C_{w,s}}{r_i \ln(r_o/r_i)} \quad (4)$$

where D_s (m²/s) is the diffusivity in the solid. In terms of diameters, and making use of the interface condition (equation 2), this can be rewritten as follows:

$$J_s = \frac{2D_s C_{w,l}}{d \ln(d_o/d)} \frac{K_s}{K_l} \quad (5)$$

Since the radial fluxes across the boundary layer (equation 1) and through the solid (equation 5) must be equal,

$$K_T (C_B - C_{w,l}) = \frac{2D_s C_{w,l}}{d \ln(d_o/d)} \frac{K_s}{K_l} \quad (6)$$

from which the surface PbLi tritium concentration $C_{w,l}$ can be written

$$C_{w,l} = \frac{C_B}{\left(1 + \frac{2\Phi_s}{K_T K_l d \ln(d_o/d)}\right)} \quad (7)$$

where the solid permeability $\Phi_s = D_s K_s$. From equations 1 and 7 an expression for the permeation flux in terms of the bulk concentration can be written:

$$J = K_T C_B \frac{\zeta}{(1 + \zeta)} \quad (8)$$

where ζ is a dimensionless number given by

$$\zeta = \frac{2\Phi_s}{K_T K_l d \ln(d_o/d)} \quad (9)$$

which has a particular significance; it indicates whether the permeation flux is limited by mass transport in the liquid or diffusion in the solid. For $\zeta \gg 1$, the flux simplifies to

$$J = K_T C_B \quad (10)$$

This is the case in which the permeation flux is limited by mass transport in the PbLi. It is identical to equation 1 with $C_{w,l} = 0$, and has no dependence on the transport properties in the solid, where in this case diffusion is fast relative to transport in the liquid PbLi. In the opposite limit $\zeta \ll 1$,

$$J = K_T C_B \zeta \quad (11)$$

or

$$J = \frac{2\Phi_s}{K_T d \ln(d_o/d)} C_B \quad (12)$$

which is the usual expression for diffusion flux through the solid, multiplied by the solubility ratio. This is the case in which transport is limited by diffusion through the solid; transport in the PbLi is comparatively fast, and thus there is no dependence on K_T .

The dimensionless number ζ is thus an important parameter in permeator design. If $\zeta \ll 1$ for a given set of design parameters, the permeation flux can be improved by choosing a higher permeability material or a thinner wall (but not by increasing the mass transport coefficient K_T); conversely, if $\zeta \gg 1$, the permeation flux can be improved by increasing K_T (but not by increasing the permeability or decreasing the wall thickness).

II.B. Axial Tritium Transport

The preceding analysis gives the radial permeation flux for a fixed bulk tritium concentration C_B . In a permeator tube, however, the bulk concentration will be continually depleted along its length due to this permeation flux. Therefore, the concentrations C_B , $C_{w,l}$, and $C_{w,s}$ will all be functions of z .

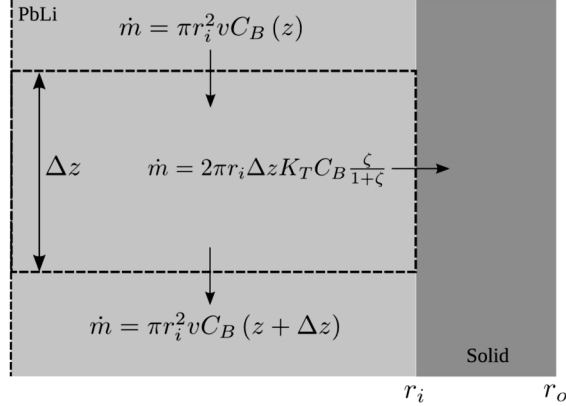


Fig. 2. Tritium balance in a PbLi fluid element.

Consider a fluid element of length Δz as shown in Figure 2. For a PbLi velocity v , the tritium flow rate (mol T/s) entering the tube is

$$\dot{m}_{in} = \frac{1}{4} \pi d^2 v C_B(z) \quad (13)$$

and the tritium flow rate exiting the element is

$$\dot{m}_{out} = \frac{1}{4} \pi d^2 v C_B(z + \Delta z) \quad (14)$$

The difference between these two is the permeation flow rate, given by equation 8 times the area:

$$\dot{m}_{perm} = \pi d \Delta z J \quad (15)$$

Balancing the fluxes gives the differential equation for $C_B(z)$,

$$\frac{\partial C_B(z)}{\partial z} = -\frac{4K_T}{vd} \frac{\zeta}{1+\zeta} C_B(z) \quad (16)$$

For an inlet tritium concentration $C_B(0) = C_0$, this is solved by

$$C_B(z) = C_0 \exp\left(-\frac{4K_T z}{vd} \frac{\zeta}{1+\zeta}\right) \quad (17)$$

II.C. Permeator Efficiency

Permeation losses (e.g. to the building and environment) elsewhere in the system are driven by the tritium concentration exiting the extraction system. The job of the permeator, then, is to remove whatever fraction of the incoming tritium necessary to achieve the desired exit concentration. This removal fraction, or efficiency, is a convenient parameter for comparison of permeator designs. Using equation 17, we can identify a permeator length L necessary to achieve a given efficiency η ,

$$C_B(L) = (1 - \eta) C_0 \quad (18)$$

which gives the following expression for the permeator efficiency:

$$\eta = 1 - \exp\left(-\frac{\tau \zeta}{1 + \zeta}\right) \quad (19)$$

where the quantity

$$\tau = \frac{4K_T L}{vd} \quad (20)$$

is another dimensionless parameter of significance: it is the ratio of time constants for axial (L/v) and radial ($d/4K_T$) transport in the permeator.

A similar expression for the permeator efficiency is given in Ref. 11,

$$\eta = 1 - \exp(A^* L) \quad (21)$$

where A^* was an unknown function of v , d , d_o , and K_T , which was later determined by fitting equation 21 to the results of numerical simulations. We see from equation

19 that for tubes (other geometries will be similar) it has a simple analytical form, which depends on not only ν , d , d_o , and K_T , but also K_l and Φ_s (see equation 9). Note that A^* must have units of inverse length.

II.D. Mass Transport Coefficient

The mass transport coefficient K_T is analogous to the empirical heat transfer coefficient, and is similarly determined by correlations of the form^{13,14}

$$\text{Sh} = \alpha \text{Re}^\beta \text{Sc}^\gamma \quad (22)$$

where the Sherwood number Sh , the analogue of the Nusselt number in heat transfer, is defined by

$$\text{Sh} = \frac{dK_T}{D_l} \quad (23)$$

and the Schmidt number Sc , the analogue of the Prandtl number, by

$$\text{Sc} = \frac{\mu}{\rho D_l} \quad (24)$$

Here D_l (m^2/s) is the diffusion coefficient in the liquid, and ρ and μ the density and viscosity, respectively. Some experimentally determined mass transport correlations applicable for turbulent flow in tubes are summarized in Table I; see also Refs. 15-17 for additional discussion.

The correlation obtained by Harriott and Hamilton¹⁸ was originally selected by Willms,⁷ and it has since been the de facto mass transport correlation for tritium in PbLi .^{3-4,11-12} That it was determined with a rather different solute and fluid (benzoic acid in glycerin water) has justifiably led to concerns about its applicability to

tritium in PbLi .¹¹⁻¹² However, the correlations in Table I, which incorporate a large number of fluids and solutes over a wide range of Reynolds and Schmidt numbers, are remarkably similar, and this gives some confidence in the applicability of the approach, at least in this geometry. In the numerical analysis of section III, we use the correlation by Linton and Sherwood,¹⁹ as the $\text{Re}^{0.83}$ dependence is verified by multiple experiments, and the $\text{Sc}^{0.33}$ dependence is thought to be more appropriate than $\text{Sc}^{0.44}$ since the latter is obtained from a comparatively narrow range of Schmidt numbers.¹⁶

But first, we retain the generic coefficients α , β , and γ in what follows so as to draw some conclusions about the permeator efficiency independent of any specific correlation. Solving equation 22 for K_T (making use of equations 23 and 24) gives

$$K_T = \frac{\alpha D_l^{1-\gamma} \nu^\beta \rho^{\beta-\gamma}}{d^{1-\beta} \mu^{a-\gamma}} \quad (25)$$

Because, regardless of the correlation used, $\beta > \gamma > 0$, all the exponents as written in equation 25 are positive.

We can substitute equation 25 into equation 19 to see how the permeator efficiency depends on the design parameters ν , d , and L . It is easiest to see if we consider the extremes of ζ ; for $\zeta \gg 1$ (liquid-limited),

$$\eta = 1 - \exp\left(-\frac{4\alpha D_l^{1-\gamma} L \rho^{\beta-\gamma}}{\nu^{1-\beta} d^{2-\beta} \mu^{\beta-\gamma}}\right) \quad (26)$$

and for $\zeta \ll 1$ (solid-limited),

$$\eta = 1 - \exp\left(-\frac{8\Phi_s L}{\nu K_l d^2 \ln(d_o/d)}\right) \quad (27)$$

TABLE I. Sherwood Number Correlation Parameters for Turbulent Flow in Tubes.

α	β	γ	Range	Reference	Notes
0.023	4/5	1/3		Ref. 13,14	Heat transfer analogy
0.023	0.83	0.44	2000 < Re < 35000 0.6 < Sc < 2.5	Ref. 20	Vaporization of nine different liquids in air
0.0328	0.77	0.33	3000 < Re < 40000 0.5 < Sc < 3	Ref. 21	Distillation of five different substances in a wetted-wall column
0.023	0.83	1/3	2000 < Re < 70000 1000 < Sc < 2260	Ref. 19	Solution of benzoic acid, cinnamic acid, and beta-naphthol in water
0.0163	0.83	0.44	Sc ~ 0.6 [†]	Ref. 22	Vaporization of water in a wetted-wall tower
0.0096	0.913	0.346	10000 < Re < 100000 430 < Sc < 100000	Ref. 18	Benzoic acid in glycerine-water, and hydroxymethylcellulose solutions

[†] This correlation has been written in the form given in Ref.16, where $\gamma = 0.44$ is assumed based on the correlation obtained in Ref. 20. As the Schmidt number was not varied in these experiments, it does not constitute an independent verification of that dependence.

In both extremes (and thus in all intermediate cases), and for all of the correlations in Table 1, the efficiency 1) increases with L , 2) decreases with v , and 3) decreases with d . The latter two points imply, perhaps counter-intuitively, that to maximize the permeator efficiency, the Reynolds number should be as low as possible (while still turbulent). There are two reasons for this; 1) increasing the diameter increases the radial distance that tritium must traverse to reach the wall, and 2) increasing the velocity does increase the radial mass transport coefficient ($\sim v^\beta$, where $\beta < 1$), but not as effectively as it increases advective transport in the axial direction ($\sim v$), i.e. the tritium is swept through the length of the permeator tube faster than it is transported radially.

III. FERRITIC STEEL PERMEATOR

Making use of the preceding analysis, we consider here the prospects of a ferritic steel vacuum permeator operating at 470 °C, a temperature limit imposed to avoid corrosion of the RAFM steel in PbLi. In order to do so, the temperature dependent material properties Φ_s , D_l , K_l , ρ , and μ must be evaluated.

III.A. RAFM Properties

A large number of diffusivity and solubility measurements in various RAFM steels have recently been compiled by Causey.²³ Based on his reported average of all these, we use here the following permeability:

$$\Phi_s = 8.72 \times 10^{-8} \exp(-41800/RT) \quad (28)$$

Here Φ_s is measured in (mol T)/(m·s·Pa^{1/2}) and $R = 8.314$ J/(mol·K).

III.B. PbLi Properties

A comprehensive summary of PbLi properties has recently been given by Mas de les Valls.²⁴ There is considerable uncertainty regarding the values of D_l and especially K_l ; upper and lower bounds for each are given in Table II with their original references.

III.B.1. The Schmidt Number

At 470 °C, and given the bounds on the diffusivity D_l in Table II,

$$40 < Sc < 150 \quad (29)$$

We note that the correlations in Table I extend far below and above this range of Schmidt numbers, further reinforcing their applicability to PbLi in these conditions.

III.C. Permeator Optimization

The preceding analysis allows for the optimization of permeator design based on system constraints. Here we solve for optimum length L and number N of tubes to minimize the total permeator volume ($N\pi d_o^2 L/4$) subject to the following constraints on tube diameter, total PbLi mass flow rate, efficiency, and allowable pressure drop:

$$d \geq 0.01 \text{ m} \quad (30)$$

$$\dot{m} = \rho N \pi r_i^2 v = 26000 \text{ kg/s} \quad (31)$$

$$\eta \geq 0.7 \quad (32)$$

$$\Delta p = f \frac{L}{d} \frac{\rho v^2}{2} \leq 1 \text{ MPa} \quad (33)$$

where f is estimated (using surface roughness $\varepsilon = 2 \text{ } \mu\text{m}$) using²⁵

$$f = -1.8 \log_{10} \left(\frac{6.9}{\text{Re}} + \left(\frac{\varepsilon/d}{3.7} \right)^{1.11} \right)^{-2} \quad (34)$$

These values are based on the ARIES-CS design⁴ and therefore may be taken as roughly applicable to a DEMO. The 70% efficiency requirement in that work was set based on the need to keep losses to the environment

TABLE II. PbLi Properties as a Function of Temperature.²⁴

Parameter	Units	Expression	Range	Reference
Density	kg/m ³	$\rho = 10520(1 - (1.13 \times 10^{-4} T))$	508-880 K	Ref. 24
Viscosity	Pa·s	$\mu = 1.87 \times 10^{-4} \exp(11640/RT)$	508-625 K	Ref. 26
Diffusivity (lower bound)	m ² /s	$D_l = 2.62 \times 10^{-9} \exp(-6630/RT)$	573-773 K	Ref. 27
Diffusivity (upper bound)	m ² /s	$D_l = 2.50 \times 10^{-7} \exp(-27000/RT)$	573-973 K	Ref. 28
Solubility (lower bound)	(mol T)/(m ³ ·Pa ^{1/2})	$K_l = 1.26 \times 10^{-3} \exp(-1350/RT)$	508-700 K	Ref. 29
Solubility (upper bound)	(mol T)/(m ³ ·Pa ^{1/2})	$K_l = 5.86 \times 10^{-2}$	573-723 K	Ref. 30

below 1 g/year (99% efficiency of room detritiation systems was also assumed). The niobium permeator in ARIES-CS (which operated at 730 °C) meets the above constraints with 6180 (2060 per 1/3 sector) tubes of 5 m in length ($v = 5$ m/s). Using the lower bound for K_l and upper bound for D_l from Table II (i.e. the most optimistic assumptions for permeator efficiency), the smallest RAFM permeator that achieves this performance at 470 °C consists of about 19400 tubes of 37.3 m in length ($v = 1.77$ m/s). This permeator is larger (in terms of volume) than a PWR steam generator and about 20 times larger than a niobium permeator that satisfies the same constraints at 730 °C (Table III). The difference in performance between these two membrane materials suggests that permeation is not strictly limited by mass transport in the PbLi, and the fact that $\zeta = 0.45$ (near unity) for the RAFM permeator confirms this. In this case, improving either the solid permeation or mass transport in the PbLi will increase the efficiency of a RAFM permeator. However, as noted previously, the only effective strategy for accomplishing the latter is to decrease the tube diameter. 1 cm was taken here as a practical lower limit, but smaller sizes can give a considerable increase in performance; it is shown in Table III that decreasing the tube diameter from 1 cm to 0.5 cm can achieve the same efficiency with almost half the size and less than half the tube length.

Using the RAFM permeator design with $d = 1$ cm described above as a base case, the uncertainty in permeator efficiency associated with both mass transport correlations and diffusivity and solubility data for PbLi is illustrated in Table IV. These are grouped into upper and

lower bounds on efficiency based on the range of diffusivity and solubility data for tritium in PbLi in Table II; the upper bound is defined by high D_l and low K_l , and the lower bound by low D_l and high K_l (see equations 26-27).

For either set of assumptions, the lowest and highest values of K_T (and by extension ζ , see equation 9) among the six correlations differ only by about a factor of two. For the upper bound assumptions, ζ is on the order of (though less than) one, and the permeator efficiency never differs from the 70% obtained above by more than 5%. For the lower bound assumptions, $\zeta \ll 1$ (i.e. permeation is solid-limited) and so the efficiency is essentially independent of the mass transport correlation.

Thus, the uncertainty in PbLi diffusivity and (especially) solubility data implies that we cannot even be sure which transport phenomenon is limiting in this case, and there is a rather large difference in the efficiencies obtained under each set of assumptions (70% and 3%). Given this fact, it would be intuitive to refer to the lower bound assumptions as conservative, but this is not necessarily the case. The most relevant metric for the extraction system is not the efficiency, but the outlet tritium concentration, which is what drives permeation losses downstream of the permeator, and it is these losses that must be minimized to maintain public and worker safety. Under the lower bound assumptions, these losses will also be lower. This tradeoff between permeation inside and outside the permeator is dependent on design details of the primary coolant system, and thus a model of this larger system is necessary to determine which set of assumptions is actually the most conservative.

TABLE III. Vacuum Permeator Comparison.

	Niobium (ARIES-CS ⁴)	RAFM (d = 1 cm)	RAFM (d = 0.5 cm)	B&W PWR steam generator ³¹
Temperature	730 °C	470 °C	470 °C	
Efficiency	0.7 (upper bound)	0.7 (upper bound)	0.7 (upper bound)	
Tubes	~6180	~19400	~80400	~15000
Tube length	5 m	37.3 m	17.0 m	20.74 m
Velocity	5 m/s	1.77 m/s	1.71 m/s	
Total Volume	3.4 m ³	69.0 m ³	38.7 m ³	61.8 m ³
ζ	124.1 (PbLi limited)	0.45 (transition)	0.43 (transition)	

TABLE IV. Sensitivity to Mass Transport Correlation (RAFM, T = 470 °C, $d = 1$ cm, $L = 37.3$ m, $v = 1.77$ m/s).

Correlation	α	β	γ	Upper Bound			Lower Bound		
				K_T (mm/s)	ζ	η	K_T (mm/s)	ζ	η
Chilton-Colburn ^{13,14}	0.023	4/5	1/3	0.32	0.64	0.65	0.14	0.025	0.029
Gilliland-Sherwood ²⁰	0.023	0.83	0.44	0.68	0.30	0.74	0.34	0.011	0.029
Johnstone-Pigford ²¹	0.0328	0.77	0.33	0.32	0.65	0.65	0.14	0.026	0.029
Linton-Sherwood ¹⁹	0.023	0.83	1/3	0.46	0.45	0.70	0.20	0.018	0.029
Kaseffian-Plank-Gerhard ²²	0.0163	0.83	0.44	0.49	0.42	0.71	0.24	0.015	0.029
Harriott-Hamilton ¹⁸	0.0096	0.913	0.346	0.54	0.38	0.72	0.24	0.015	0.029

IV. CONCLUSIONS

There are a number of important points to be emphasized based on the preceding analysis. By deriving an analytical solution, we have shown that the permeator efficiency (i.e. tritium removal fraction) can be described with a simple function of two dimensionless numbers: one indicating whether radial transport in the liquid PbLi or solid membrane is limiting, and another that is a ratio of axial to radial transport times. Based on the form of correlations for the mass transport coefficient, the permeator efficiency is increased by increasing the tube length, decreasing the tube diameter, and decreasing the PbLi velocity in the tube (e.g. by adding more tubes).

A review of numerous experimentally measured mass transport coefficients finds that these are quite consistent with each other and with the expected form based on the heat transfer analogy, and that they span a wide range of Schmidt numbers, including those relevant to PbLi. Therefore, there is no reason to suspect they are not applicable to tritium transport in PbLi, nor does the selection of mass transport correlation represent a large source of uncertainty in the present analysis relative to other parameters.

The PbLi diffusivity and solubility (especially the latter), on the other hand, are considerable sources of uncertainty. Considering upper and lower bounds for D_l and K_l from the literature, for the RAFM permeator parameters determined here, it can only be said that efficiency is between 3% and 70%.

On the last point, we again emphasize the fact that there is a tradeoff here: high D_l and low K_l increase permeation everywhere in the system, both in the permeator (where it is desired) and elsewhere (where it is not). Taking a lower D_l and higher K_l implies a less efficient permeator, but also lower losses elsewhere in the system, i.e. in this case less efficiency is needed. This is a point we intend to investigate further in future work.

Regardless of which assumptions are made regarding these PbLi properties, it is clear that a permeator with RAFM membranes must be significantly larger (~20 times the volume) than previous concepts based on high temperature group 5 metals. In light of that fact, future efforts to identify higher permeability (and/or high-temperature compatible) materials that meet necessary corrosion, activation, and oxidation constraints are probably warranted.

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