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Changing the World's Energy Future

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Phase-field simulations of fission gas bubbles in high burnup UO_2 during steady-state and LOCA transient conditions

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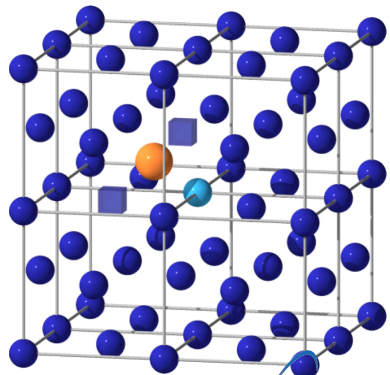
Idaho National Laboratory

Michael Cooper, Topher Matthews

Los Alamos National Laboratory

Multiscale Modeling for performance simulation

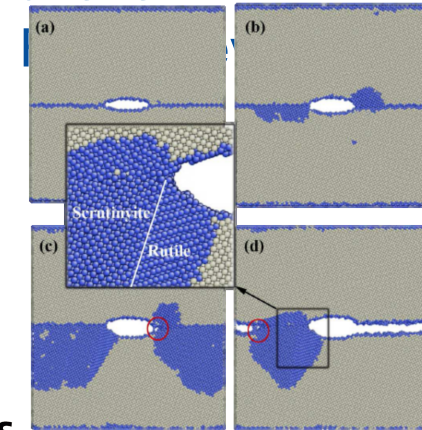
- BISON: Fuel performance code developed at INL
- Inform BISON with atomistic and mesoscale simulations
- Click-Marmot: MOOSE-based phase-field simulation code
 - Second level
 - Third level



Los Alamos
NATIONAL LABORATORY
EST. 1943

nanometers First Principles

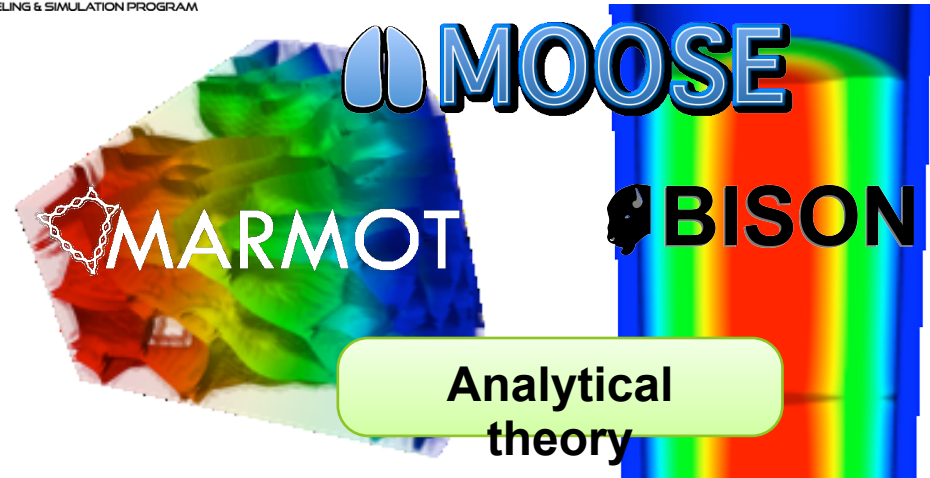
- Identify critical bulk mechanisms
- Determine bulk properties



100's of nanometers Molecular Dynamics

- Identify interfacial mechanisms
- Determine interfacial properties

NEAMS
NUCLEAR ENERGY ADVANCED MODELING & SIMULATION PROGRAM



microns Mesoscale

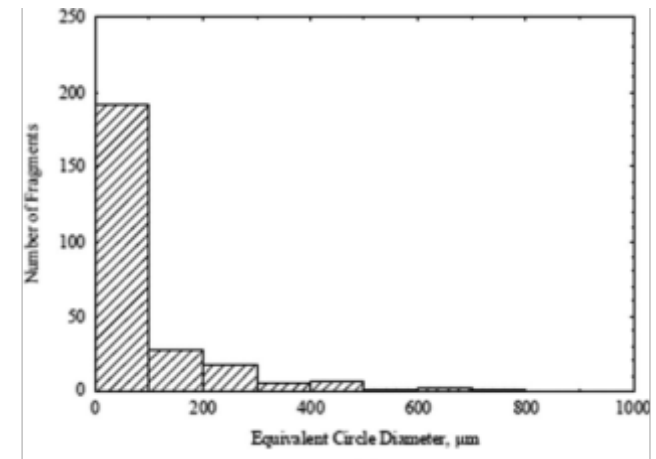
- Predict microstructure evolution
- Determine impact on properties

millimeters and up Engineering Scale

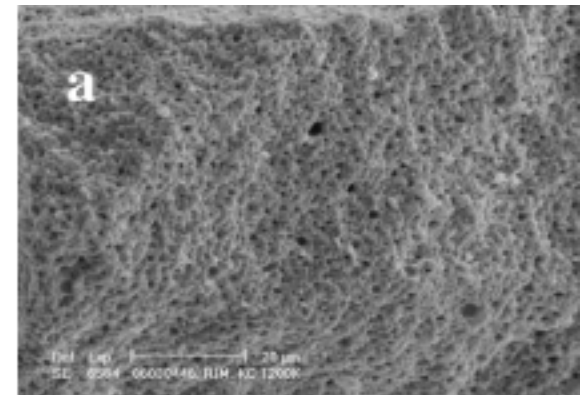
- Use analytical theory
- Predict fuel performance

Linkage from Material/Particle Verization of high-burnup UO_2

- Potential to occur during loss-of-coolant accident (LOCA)-type temperature transients
 - Second level
- Formation of fine fragments <100 micron in size
 - Fourth level
- Fine particles can potentially escape into coolant from burst cladding during LOCA
 - Fifth level
- Industry would like to understand this problem better to strengthen their case for increasing fuel burnup limits
- Hypothesized mechanism: During LOCA, trapped gas in bubbles heats up and becomes overpressurized; cracking initiates at these overpressurized bubbles



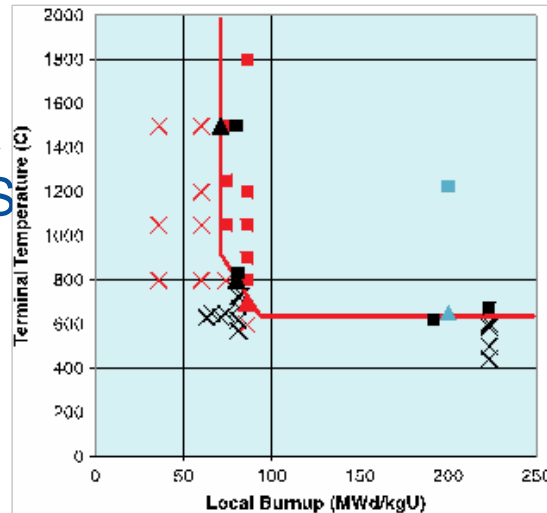
Turnbull et al., Nuc. Sci. & Eng., 179,477 (2015).



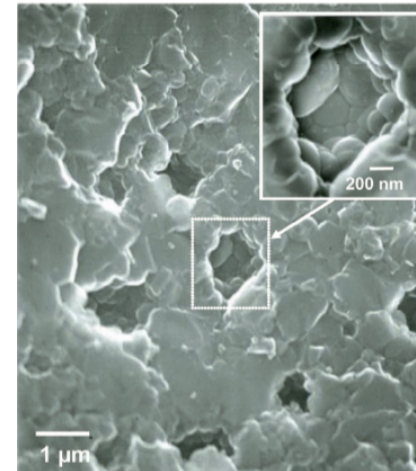
Hiernaut et al, JNM 377, 313 (2008).

BISON Nondestructive Fuel Characterization

- Click – S



Turnbull et al., Nuc. Sci. & Eng., 179,477 (2015).

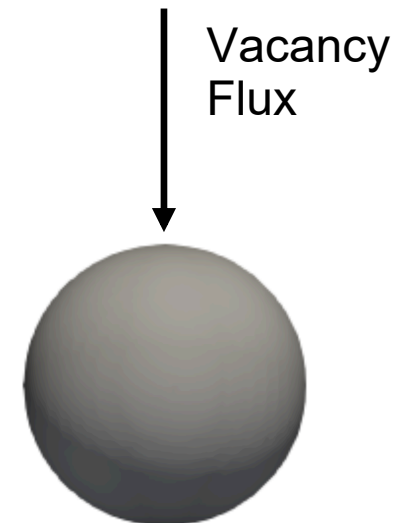


High-burnup structure in UO₂
[Sonoda et al., NIMB, 2002].

- Current model is empirical, based on burnup and temperature
- Pulverization is predominantly observed to occur in regions where high-burnup structure (HBS) has partially or completely formed
 - HBS: Grain size decreases to ~150–200 nm, micron-sized bubbles form with multiple grains intersecting each bubble
- Goal: Develop a physics-based criterion for pulverization in BISON that accounts for microstructure
 - Focus on HBS

BSK to bubble response to LOCA transient

- Bubbles in HBS region are $\sim 1 \mu\text{m}$ and believed to be overpressurized relative to equilibrium (based on observed dislocation punching around bubbles):
 - Second level
 - Third level
 - Fourth level
- Overpressurized bubbles exert compressive stress in the radial direction on the surrounding matrix.
 - Fifth level
- During LOCA transient, temperature and therefore bubble pressure increases further, causing stress in the matrix to increase further. Compressive stress leads to increased vacancy flux to bubble, causing bubble growth.
- **Key Questions:**
 - Does significant bubble growth occur during duration of a LOCA transient?
 - What is the pressure response to a given temperature transient?



Phase-field model for essential physics

- Single order parameter η to represent gas bubble and solid matrix phase
 - Current model does not consider grain boundaries
 - Fourth level
- Track vacancies and fission product gas atoms
 - Use Xe properties for fission product gases
 - Source terms for production, sink term to limit vacancy concentration to steady-state
- Chemical and elastic energy contributions
- Solid-bubble interfacial energy
 - Kim-Kim-Suzuki (KKS) approach to remove bulk energy contribution to interfacial energy
- Surface tension of bubble-matrix interface
- Xe gas pressure

$$\eta = 0$$



$$\eta = 1$$

Phase-field model free energy functional

$$F = \int_V \left[f_{chem} + Wg(\eta) + \frac{\kappa}{2} |\nabla \eta|^2 + f_{el} \right] dV$$

- Click to edit text

- **Second level**
bulk chemical free energy density. $h(\eta)$ is a smooth interpolation function.

$$f_{chem} = [1 - h(\eta)] f_{chem}^m(c_a^m, c_g^m) + h(\eta) f_{chem}^b(c_a^b, c_g^b)$$

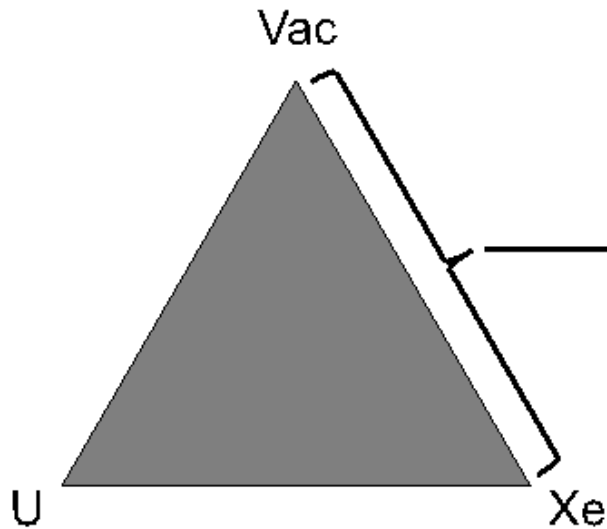
- **Fourth level**
• **Fifth level**
chemical free energy of the matrix phase. Fit a parabolic approximation to ideal solution energetics:

$$f_{chem}^m \approx f_{chem}^m - \frac{k_a^m}{2} (c_a^m - c_a^{m,ref})^2 + \frac{k_g^m}{2} (c_g^m - c_g^{m,ref})^2$$

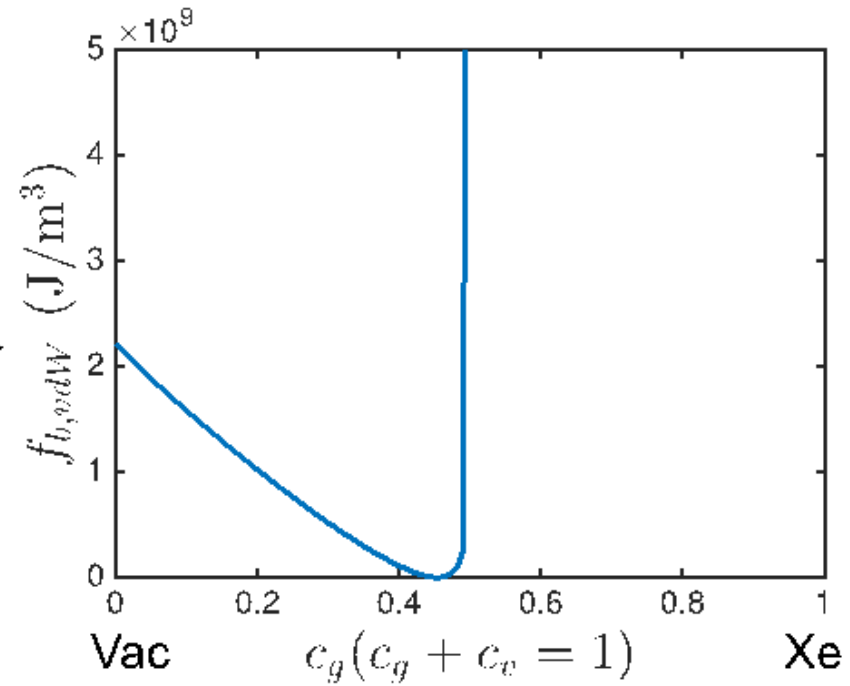
- chemical free energy of the bubble phase. The bubble is considered to be a mixture of gas atoms and U-site vacancies. Energy given by the Helmholtz free energy of a van der Waals gas:

$$f_{chem}^b = c_g^b \frac{kT}{V_{\alpha}} \left[\ln \left(\frac{1}{n_{Q2} \left(\frac{V}{V_0} - b \right)} - 1 \right) \right] + \frac{k_g}{2} (1 - c_g^b - c_v^b)^2 + f_0$$

Equilibrium state for the gas phase



Gibbs triangle: U lattice sites



Helmholtz free energy density: Van der Waals gas

Phase field model Elastic energy

- **Computing elastic energies and stresses (Voigt-Taylor scheme):**

– Second level

$$f_{el} = [1 - h(\eta)] f_{el}^a + h(\eta) f_{el}^b \quad f_{el}^a = \frac{1}{2} C_{ijkl}^a \epsilon_{ij}^{el,a} \epsilon_{kl}^{el,a} \quad f_{el}^b = \frac{1}{2} C_{ijkl}^b \epsilon_{ij}^{el,b} \epsilon_{kl}^{el,b}$$

• Third level

- Mechanical equilibrium equation:

– Fourth level

$$\nabla \cdot \sigma_{ij} = \nabla \cdot [(1 - h(\eta)) \sigma_{ij}^a + h(\eta) \sigma_{ij}^b] + \sigma_{ij}^v = 0 \quad \sigma_{ij}^a = C_{ijkl}^a \epsilon_{kl}^{el,a}$$

• Fifth level

- Eigenstrain due to vacancies:

$$\epsilon_{ij}^{el,v} = \epsilon_{ij} - \epsilon_{ij}^0 = \epsilon_{ij} - (c_v - c_v^0) \epsilon_1^0 \delta_{ij}$$

- Bubble pressure:

$$\sigma_{ij}^b = - \left(\frac{RT}{\frac{V_s}{c_v^0} b} \right) \mathbf{I} + C_{ijkl}^b \epsilon_{kl}^{el,b}$$

- Surface tension:

$$\sigma_{ij}^s = - \left[W g(\eta) + \frac{\kappa}{2} |\nabla \eta|^2 \right] \mathbf{I} - \kappa \nabla \eta \otimes \nabla \eta$$

Evolutionary Modeling

- Allen-Cahn for order parameter:

Second level

Third level

$$-L \left[\frac{db}{da} \left(\frac{d\eta}{dt} - \mu_V \right) - \mu_V (c_V^{\infty} - c_V^0) - \mu_g (c_g^{\infty} - c_g^0) \right] - W \frac{d\eta}{dt} - \kappa \nabla^2 \eta$$

Fifth level

- Cahn-Hilliard for vacancy and gas concentration (source for vacancies and gas atoms, sink for vacancies to approximate recombination):

$$\frac{\partial c_v}{\partial t} = \nabla \cdot (M_v \nabla \mu_v) - s_v - K_r c_v^{\infty}$$

$$\frac{\partial c_g}{\partial t} = \nabla \cdot (M_g \nabla \mu_g) + s_g$$

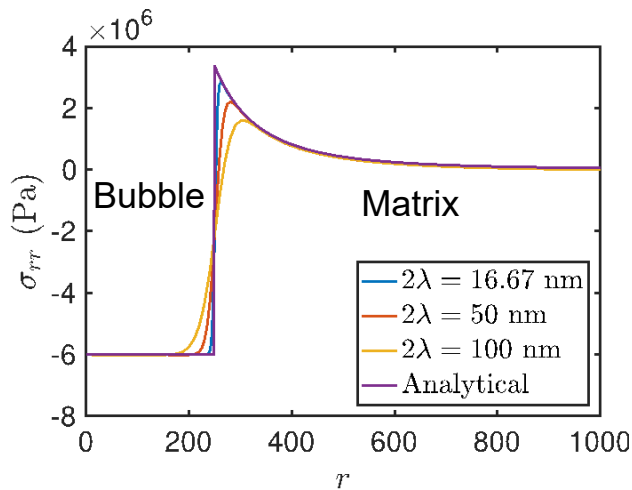
- Mobilities are a function of defect diffusivities:

$$M_v = \frac{bD_v^b + (1-b)D_v^m}{kT/c_v} \quad M_g = \frac{bD_g^b + (1-b)D_g^m}{kT/c_g}$$

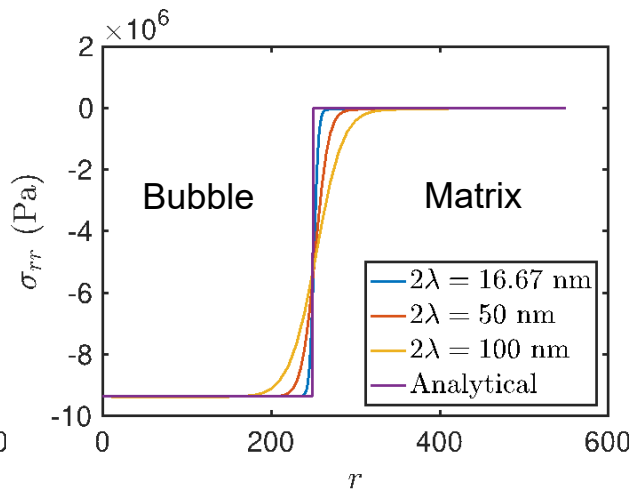
- + KKS system constraints

Phase-field Model verification and testing

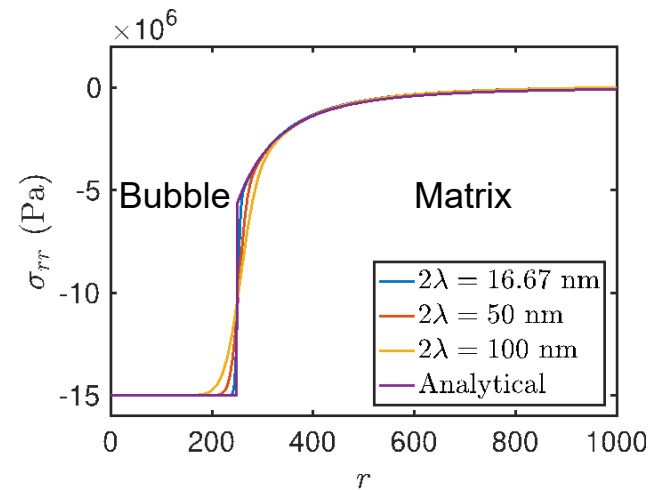
- Considered stress state in equilibrium, underpressurized, overpressurized bubbles (1D simulation in radial coordinates)
- Click to edit text
 - Equilibrium; $\sigma_{rr} = 0$ in surrounding solid matrix
 - Under/overpressurized: $\sigma_{rr} = +/-$, corresponding to tensile/compressive stress in surrounding matrix
 - Converge to analytical solution as interface width (λ) decreases



Underpressurized

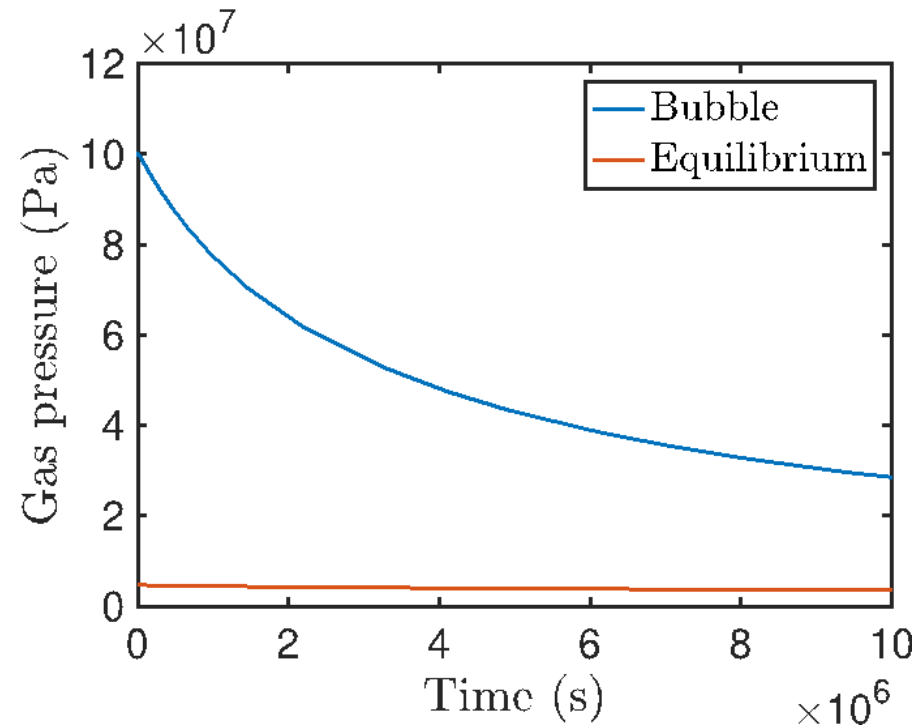
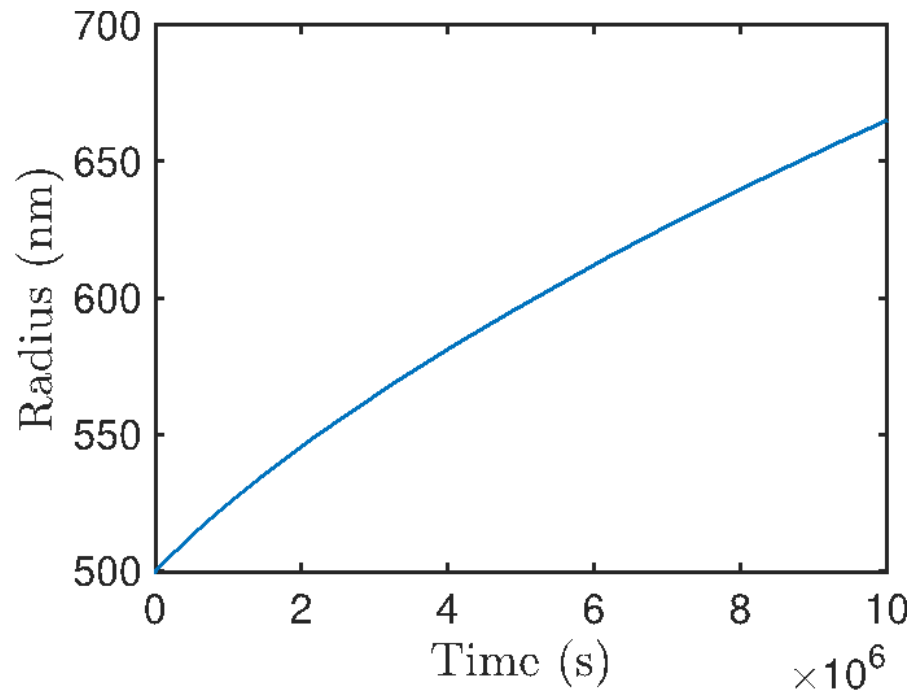


Equilibrium



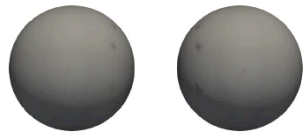
Overpressurized

Bubble growth during steady-state operation

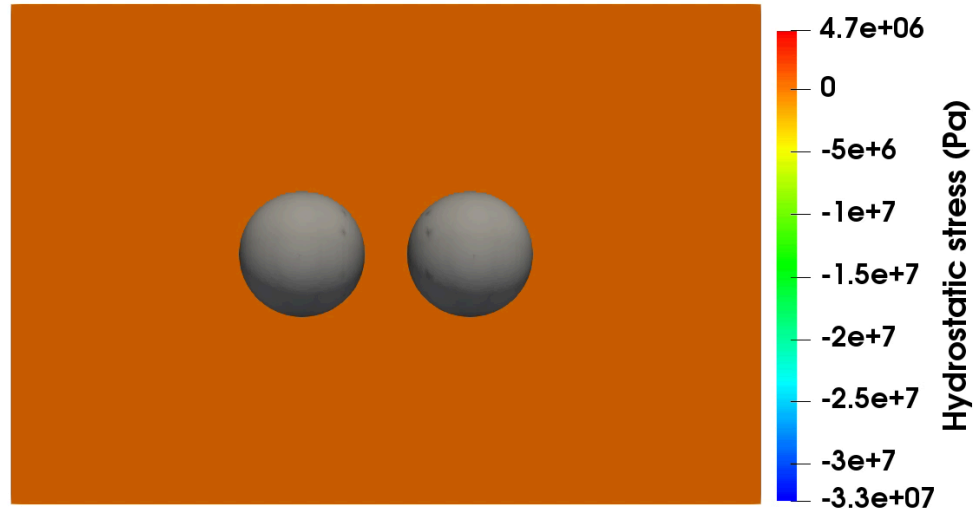


- Assume bubble pressure is 100 MPa in initial conditions
 - Upper bound based on dislocation punching pressure
- Bubble pressure decreases during growth but remains well above equilibrium pressure
 - Increased likelihood of fragmentation during LOCA

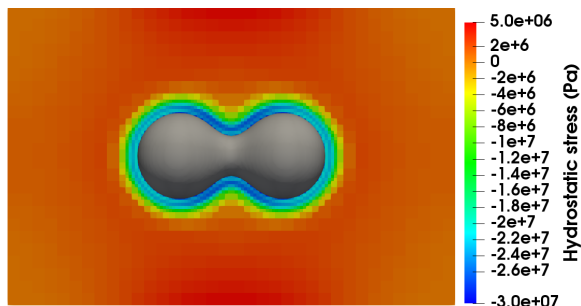
Bubble growth in a steady-state operation



rel



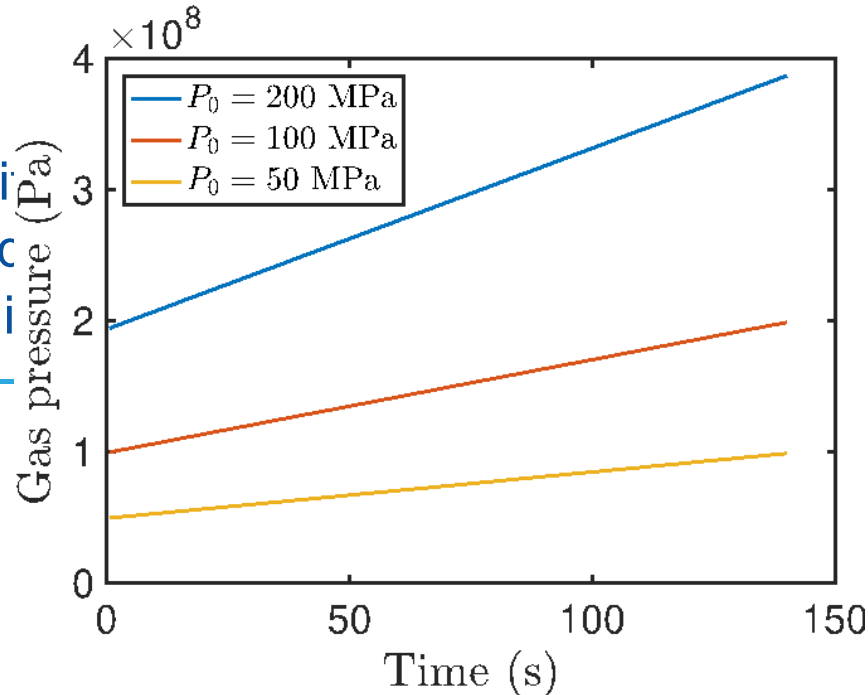
- 3D simulation to 1.5×10^7 s, 2 bubble impingement, initial radii of 300 nm
- Hydrostatic stress surrounding bubbles
 - Region of enhanced compressive hydrostatic stress in “neck”



$t = 1.3 \times 10^6$ s

Click to respond Mastering the LOCA transient

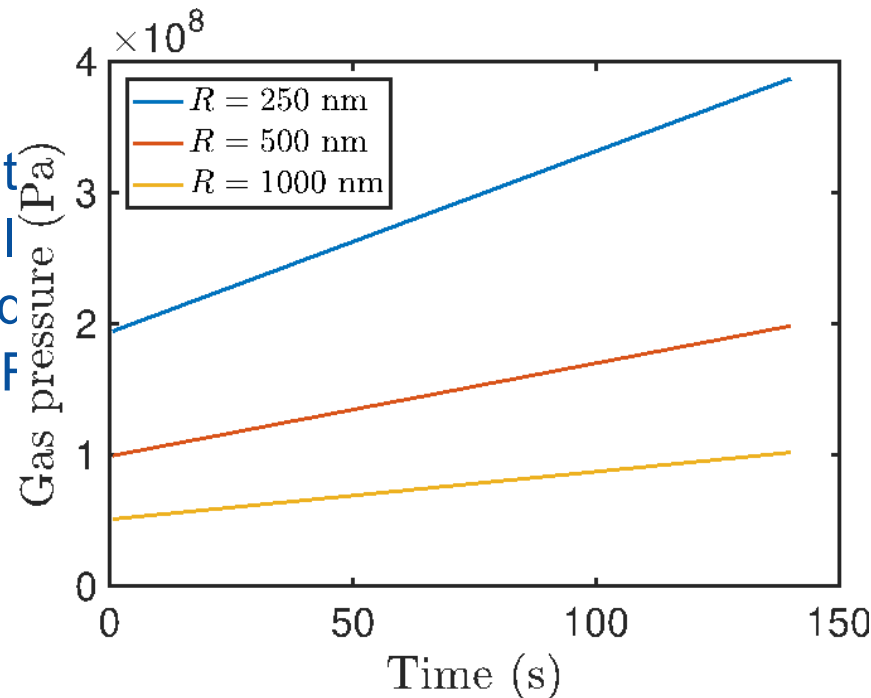
- Click to edit
 - Second
 - This



- Temperature ramp from 700 to 1400 K at 5 K/s
- Consider fixed bubble radius of 250 nm in initial conditions
- Maximum initial pressure set to $P_0 = 200$ MPa (upper bound based on dislocation punching); vary P_0 for fixed bubble size
- No significant change in bubble radius for each case

Click to respond Mastering the LOCA transient

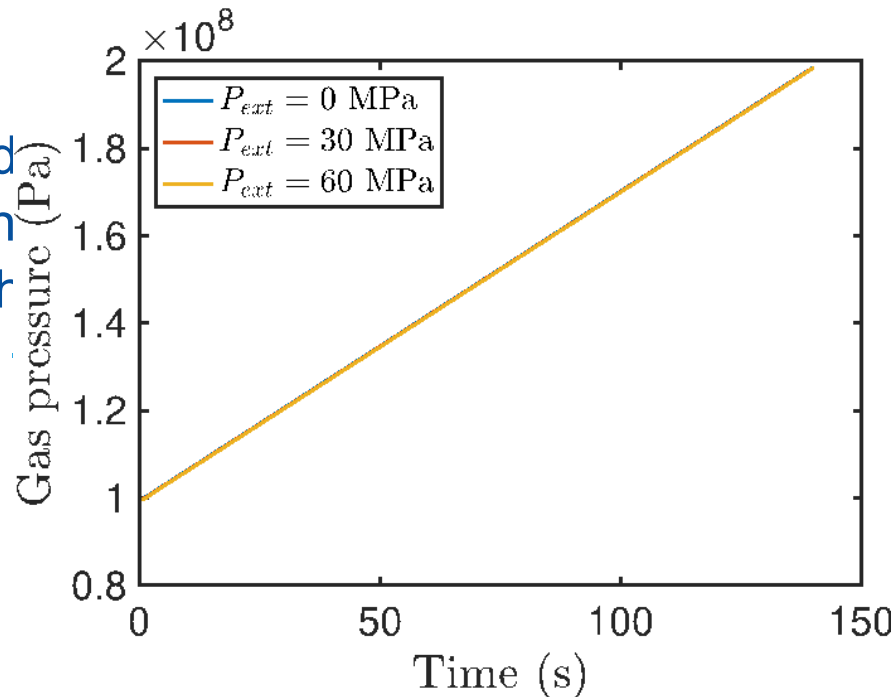
- Click to edit t
 - Second l
 - Thirc
 - F



- Vary initial radius: 250 nm, 500 nm, 1000 nm
 - Change domain size to maintain 10% porosity
- Initial pressures set at upper bound estimate from dislocation punching: 200, 100, 50 MPa.
- No significant change in bubble size

Click to respond Mastering the LOCA transient

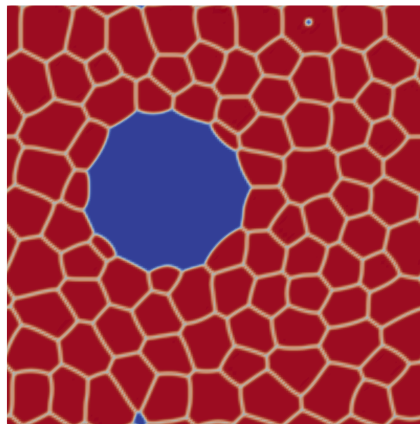
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- Second
- Third



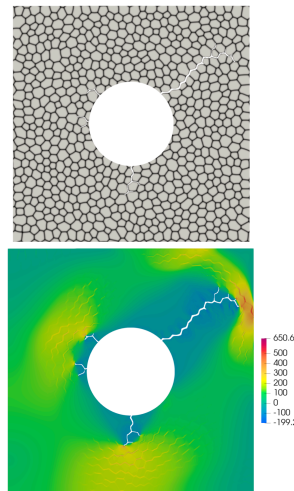
- Vary external pressure at simulation domain boundary, P_{ext} , for constant bubble $R = 500$ and $P_0 = 100$ MPa
- No significant size change; pressure transient unchanged
- P_{ext} may have a stronger impact on fracture behavior
- **Based on these results, do not need to consider bubble size change in phase-field fracture model**

Mechanistic Model for pulverization of high burnup UO_2 fuel during LOCA

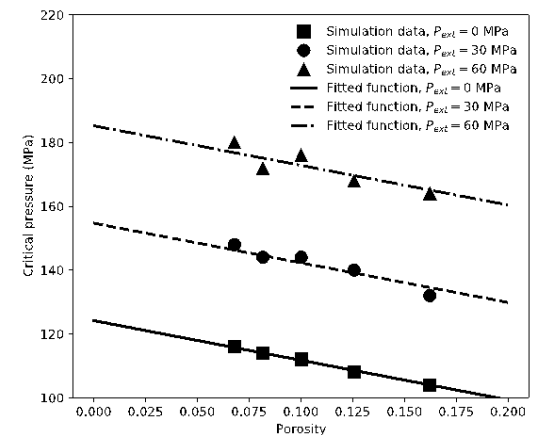
- Estimate LOCA HBS bubble pressure using grand potential-based multi-grain phase-field model
 - Second level
 - Initial pressure input to BISON
- Phase-field fracture level simulations
 - Fit to get BISON criterion as a function of porosity, external stress
- BISON model validation in progress



Bubble
Evolution
& Pressure



Fracture during LOCA



Bison Criterion

Conclusions: Major findings up UO₂ response to LOCA transients

- Developed new phase-field model that accounts for effects of surface tension and gas bubble pressure to understand non-equilibrium bubbles
 - Second level
 - Third level
- Bubble size did not change significantly during LOCA transients
 - Fourth level
- Pressure as a function of time determined for given transients
 - Lack of bubble size change allows linear pressure increase with temperature in phase-field fracture model
- Initial implementation of BISON pulverization model completed, validation in progress
 - Fifth level

Thank you!

**Funding Support: DOE-NE
NEAMS Program**



Questions?